CHARACTER OF QUASI-BANDS IN ¹⁵⁰Sm USING IBM

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Abstract: The interacting boson model-1 Hamiltonian is used to describe the energy spectrum for six quasi-bands in ¹⁵⁰Sm. In this work, g-, β -, and γ - bands in ¹⁵⁰Sm isotope are studied by using interacting boson model-1. It is found that the calculated energy values and B(E2) values have good agreement with experimental data. The calculated energy values and B(E2) values are also compared with the Dynamic Pairing-Plus-Quadrupole (DPPQ) Model.

1. Introduction

The interacting boson model-1 (IBM-1) of Arima and Iachello [1] has been successful in describing the collective nuclear properties in the medium mass nuclei. Earlier systematic studies of ¹⁴⁶⁻¹⁵⁴Sm have been performed using interacting boson model [2] and Dynamic Pairing-Plus-Quadrupole (DPPQ) Model [3,4,5]. Very recently Diab [6] presented electrical monopole transition structure of ¹⁵⁰Sm isotope. Simpson et al. [7] successfully interpreted the isomers of ^{156,158}Sm using Quasi-particle Rotor Model. Sharma and Kumar [8] presented a fresh analysis of g- , β -, and γ - bands in ¹⁵⁰Sm using IBM. This search is now extended to calculate B(E2) values for (g --> g), (β --> β), (γ --> g) and (γ --> β) transitions using IBM and compare with experimental data [9,10]. The study of quasi-bands in ¹⁵⁰Sm has not been studied sufficiently, thus the present study is interesting to investigate. Recent experimental data of energies have been taken from the website of National Nuclear Data center, Brookhaven National Laboratory, USA [16].

2. The Interacting Boson Model

The phenomenological Interacting Boson Model -1 (IBM-1) initially introduced by Arima and Iachello [1] has been rather successful in describing the collective properties of several medium and heavy mass nuclei. In the first approximation, only pairs with angular momentum L = 0 (called s-bosons) and L = 2 (called d-bosons) are considered. The model has associated with it an inherent group structure, which allows for the introduction of limiting symmetries called SU(5), SU(3) and O(6). However, in a more general case the full IBM-1 Hamiltonian has to be used, which has several forms [1]. The multi-pole form of the interacting boson model-1 Hamiltonian is given by

$$H = \varepsilon \hat{n}_{d} + a_{0} \left(\hat{P}^{+} \cdot \hat{P} \right) + a_{1} \left(\hat{L} \cdot \hat{L} \right) + a_{2} \left(\hat{Q} \cdot \hat{Q} \right) + a_{3} \left(\hat{T}_{3} \cdot \hat{T}_{3} \right) + a_{4} \left(\hat{T}_{4} \cdot \hat{T}_{4} \right),$$

where the interaction parameters in the PHINT Program are given below: $\varepsilon = \text{EPS}$, $a_0 = 2\text{PAIR}$, $a_1 = \text{ELL}/2$, $a_2 = \text{QQ}/2$, $a_3 = 5\text{OCT}$ and $a_4 = 5\text{HEX}$.

2.1. E2 Transitions

The E2 transitions provide more stringent test of the model. The general E2 transition operator is given by

$$T^{(E2)} = \alpha_2 \left[d_+ \hat{S} + S^+ \hat{d} \right]^{(2)} + \beta_2 \left[d_+ \hat{d} \right]^{(2)}.$$

The coefficient α_2 called the boson effective charge is an overall scaling factor for all B(E2) values which is determined from the fit to the $B(E2, 2_1^+ - 0_1^+)$ value. The coefficient β_2 may be determined from the quadrupole moment $Q(2_1^+)$. The ratio $\beta_2/\alpha_2 = \chi = -1.32$ in the SU(3) limit and is reduced to zero in the O(6) limit. In the "FBEM" program the corresponding parameters are $\alpha_2 = E2SD$, $\beta_2 = (1/\sqrt{5})(E2DD)$.

3. Results and Discussion

3.1 Energy Spectrum

The triplet of 4_1^+ , 2_2^+ and 0_2^+ states is well broken in ¹⁵⁰Sm. The 'H' Hamiltonian is used to calculate the spectrum of the ¹⁵⁰Sm. In ¹⁵⁰Sm isotope, the six positive parity bands, i.e. g-, β - and γ -, β_2 -, β_3 - and γ_2 -bands are well established [15]. Sakai [15] listed the g-band up to 14^+ , β_1 -band up to 4^+ , γ_1 - band up to 5^+ , β_2 - band up to 4^+ , γ_2 -band up to 3^+ and β_3 - band up to 2^+ states in ¹⁵⁰Sm isotope. The results of present calculation for energy spectrum are presented in Table 1 and there is a good agreement between theory and experiment (see Figure 1).



Fig. 1. Comparison of experimental and calculated (IBM) energy spectra of ¹⁵⁰Sm.

State	Expt.[15,16]	Present work	DPPQ[5]	
2g	0.33395	0.4281	0.338	
4g	0.77335	0.8612	0.747	
6g	1.27885	1.3195		
8g	1.8371	1.8177		
$0\beta_1$	0.74042	0.6640	1.049	
$2\beta_1$	1.04614	0.9095	1.460	
$4\beta_1$	1.44917	1.3451		
$2\gamma_1$	1.19381	1.1439	1.782	
$3\gamma_1$	1.50453	1.4936		
$4\gamma_1$	1.64260	1.6544		
$5\gamma_1$	2.0204	2.0068		
$0\beta_2$	1.2555	1.2277		
$2\beta_2$	1.4173	1.5549		
$4\beta_2$	(1.8194)	1.9605		
$2\gamma_2$	1.7938	1.8226		
$3\gamma_2$	2.0633	2.1632		
0β3	(1.7610)	1.6650		
2β ₃	2.0055	2.1557		

Table 1. The values of energy (in MeV) for ¹⁵⁰Sm.

3.2 B(E2) Values

The absolute B(E2) values are presented in table 2, where the experimental data [9,11] is compared with the present calculation and other previous work [5,9,11,12,13]. The B(E2) values are available for 15 transitions. In the present calculation the B(E2) values for (g -> g) transition increases with spin I⁺. The same feature is also observed by Yen et al [14], however, in the experiment B(E2) values increases with spin up to 4⁺ and decreases when spin increases from 8⁺ to 10^+ . For ($\beta \rightarrow g$), ($\beta \rightarrow \beta$), ($\gamma \rightarrow g$) and ($\gamma \rightarrow \beta$) transitions the theoretical values lie very close to the experimental data. In the effective IBM-1 calculation [13] the B(E2; 2g --> 0g) value is 8 times smaller than observed value, for (2 $\beta \rightarrow 0g$) transition it gives 135 times larger value than observed and for ($2\gamma \rightarrow 2g$) transition the calculated value is 20 times smaller than experiment, but in present calculation all these three B(E2) values lie very close to the observed values. Yen et al. [14] used the effective IBM calculation and included the Z = 64 proton sub-shell effect for N< 90 region. They [12] used two sets of calculation, i.e. MI and MII, in MI the proton boson were counted from Z = 50closed shell, however, in MII set the Z = 64 sub-shell was included and obtained better results from MII [12]. However, in the present work proton bosons are counted from Z = 50 and we have obtained better results than MII set of Han et al. [12].

I i -> I f	Expt.	Present	D	DPPQ		IBM ^f		IBM ^g
	b		d	e	с	MI	MII	
2g>0g	0.264(12) ^a	0.269	0.41	0.186	0.275	0.355	0.318	0.034
4g>2g	0.49(4)	0.5339	0.73	0.186	0.51	0.708	0.566	0.54
6g>4g		0.7465						
8g>6g	0.435 ^c	0.8738				0.934	0.609	
10g>8g	0.447 ^c	0.9076				0.700	0.512	
4g>2β	0.0106 ^c	0.1201			0.139	0.08	0.05	0.006
0β>2g	0.218(26) 0.26(3) ^c	0.686	0.47	0.43	0.42			
2β>0g	0.004(2) ^c	0.0146	0.008	0.0074	0.02	0.09	0.021	0.54
2β>2g	0.27(15) 0.043(20) ^c	0266	0.12	0.108	0.181	0.14	0.11	
2β>4g	0.55(30) 0.17(10) ^c	0.216	0.10	0.10	0.077			
4β>2g		0.013		0.0004	0.007			
4β>4g		0.127		0.063	0.105			
2β>0β	0.56(31)	0.23	033					
2γ>0g	$0.009(4) \\ 0.009(2)^{c}$	0.004		0.015	0.02	0.0001	0.004	0.001
2γ>2g	0.29(11) 0.039(14) ^c	0.082	0.029	0.027	0.024	0.005	0.08	0.002
2γ>4g	0.028(13) $0.019(10)^{c}$	0.129	0.059	0.054	0.087	0.036	0.025	
2γ>0β	0.034(16)	0.1297	0.097					
2γ>2β	0.4(3)	0.3323	0.55					
4γ>2g		0.0042		0.009	0.022			
4γ>4g		0.0677		0.035	0.032			
Reference [9] Reference [5]	^b Re	eference [10] eference [12]	^c R ^g R	^c Reference [11] ^g Reference [13]		^d Reference [4]		

Table 2. Absolute B(E2; I i -M f) values (in e2 b2 unit) for 150Sm.

4. Conclusion

In the present work we studied systematically the lower and higher states of lower and higher bands, absolute B(E2) values of ¹⁵⁰Sm, which lie near the O(6) limit of IBM-1[1] because $R_4 = 2.31$. In ¹⁵⁰Sm, the triplet of 4_1^+ , 2_2^+ and 0_2^+ is well broken, which have a splitting of 0.4534 MeV and the value of asymmetry parameter $\gamma_0 = 20.41^\circ$. The energy difference $\Delta E = \{(E_{2g} + E_{2\gamma}) - E_{3\gamma}\} = 7.0$, which

indicates ¹⁵⁰Sm isotopes are not rotational nuclei because for rotational nuclei $\Delta E = 0$. Thus, present calculations reflects that ¹⁵⁰Sm isotope lies in the transition from SU(5) to SU(3) limit of IBM-1[1].

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