

G. A. ADOMIAN

OBTAINING FIRST AND SECOND ORDER STATISTICS $\langle y \rangle$ and $R_y(t_1, t_2)$ IN STOCHASTIC DIFFERENTIAL EQUATIONS FOR THE NONLINEAR CASE

Previous papers [1, 2] by the author have recently developed a convergent method of successive approximation for determination of the first and second order statistics, $\langle y \rangle$ and $R_y(t_1, t_2)$ for the solution of the nonhomogeneous nonlinear stochastic differential equation

$$Ly + n(y, \dot{y}, \dots) = x(t) \quad (1)$$

where $x(t, \omega)$, $t \in T$, $\omega \in (\Omega, \mathcal{F}, \mu)$, a probability space, is a stochastic process; L is a linear stochastic operator [3, 4, 5] given by $L = \sum_{v=0}^n a_v(t, \omega) d^v/dt^v$, an n th order differential operator with stochastic coefficients $a_v(t, \omega)$, $t \in T$,

$$\omega \in (\Omega, \mathcal{F}, \mu), \text{ for } v = 0, 1, \dots, n-1 \text{ and } n(y, \dot{y}, \dots) =$$

$$= \sum_{\mu=0}^{m_m} b_\mu(t, \omega) (y^{(\mu)})^{m_\mu}$$

representing a nonlinear operator with stochastic coefficients b_μ on $T \times \Omega$, where $y^{(\mu)}$ of course represents the μ th derivative. The processes may be complex and defined on different probability spaces as well but this will not be shown here.

We assume that the operator L is separable into a sum of deterministic and random parts, L and R respectively, i. e., $L = L + R$. Assuming $a_n > 0$ and is deterministic, while a_v , for $v = 0, 1, \dots, n-1$ may be stochastic, $a_v = \langle a_v \rangle + a_v(t, \omega)$ where a_v is a zero-mean random process, thus $L = \sum_{v=0}^n \langle a_v \rangle d^v/dt^v$ and $R = \sum_{v=0}^{n-1} a_v(t, \omega) d^v/dt^v$. We assume that $\langle a_v \rangle$ exists and is continuous on T . The development also assumes that L is invertible, i. e., a Green's function $e(t, \tau)$ exists. Further, $\langle b_\mu \rangle$ exists and the random or fluctuating parts of a_v and b_μ , given by a_v and β_μ , are continuous, a. e. We assume the input x is bounded on T and that derivatives of a_v and b_μ to the necessary order are bounded, and finally that the input process and coefficient processes are statistically independent. Now we write (1) as

$$y(t, \omega) = L^{-1} x(t, \omega) - L^{-1} R y(t, \omega) - L^{-1} n(y, \dot{y}, \dots) \quad (2)$$

or in terms of the Green's function

$$y(t, \omega) = \int_0^t e(t, \tau) \times (\tau, \omega) d\tau - \int_0^t e(t, \tau) \sum_{\nu=0}^{n-1} z_\nu(\tau, \omega) (d^\nu y(\tau, \omega)/d\tau^\nu) d\tau - \\ - \int_0^t e(t, \tau) \sum_{\mu=0}^{m-1} b_\mu(\tau, \omega) (d^\mu y(\tau, \omega)/d\tau^\mu)^{m,1} d\tau. \quad (3)$$

For the case of zero initial conditions [1, 3], Adomian's "stochastic Green's theorem" allows writing the second term in terms of an adjoint operator as

$$\int_0^t R_1^\dagger [e(t, \tau)] y(\tau, \omega) d\tau$$

where

$$R_1^\dagger [e(t, \tau)] = \sum_{k=0}^{n-1} (-1)^k (d^k/d\tau^k) [z_k(\tau) e(t, \tau)] = K(t, \tau, \omega).$$

This term has also been given elsewhere in terms of a "stochastic resolvent kernel" by Sibul [6] and by Adomian [3]. Letting

$$y = \sum_{i=0}^{\infty} (-i)^i y_i \text{ where } y_0 = F(t) = L^{-1} x(t) + \Phi \text{ where } \Phi$$

satisfies the homogeneous equation $Ly = 0$, and expanding n in a Taylor series, we obtain an iterative solution [1] with each term y_i for $i \geq 1$ depending on the *previous* y_{i-1} . Since we eventually work back to y_0 , ensemble averages will separate because of the statistical independence of x from the coefficient processes without closure approximations or white noise assumptions.

The objective of this paper is to show further details of the determination of the mean solution $\langle y \rangle$ and the twopoint correlation function $R_y(t_1, t_2)$, or the covariance $K_y(t_1, t_2)$, for $y(t, \omega)$. To obtain these statistical measures [3], y must, of course, be averaged over the space Ω . (If the probability space are different for z_ν , β_μ , x , then the averaging must be carried out over the appropriate spaces). Similarly to obtain $R_y(t_1, t_2)$, we find $\langle y(t_1) y(t_2) \rangle$ or $\langle y(t_1) \bar{y}(t_2) \rangle$ in the complex case. To correspond to earlier notation from the linear case, assume for the moment that $n = 0$ and in terms of the stochastic resolvent kernel Γ

$$y(t, \omega) = F(t, \omega) + \int_0^t K(t, \tau, \omega) y(\tau, \omega) d\tau = \\ = F(t, \omega) + \int_0^t \Gamma(t, \tau, \omega) F(\tau, \omega) d\tau$$

where

$$\Gamma(t, \tau, \omega) = \sum_{m=1}^{\infty} (-1)^{m-1} K_m(t, \tau, \omega),$$

$$K_m(t, \tau, \omega) = \int_0^t K(t, \tau_1, \omega) K_{m-1}(\tau_1, \tau, \omega) d\tau_1,$$

$$K_1 = K.$$

Now including the nonlinear term

$$\begin{aligned} \langle y \rangle &= \langle F(t) \rangle - \int_0^t \langle \Gamma(t, \tau) \rangle \langle F(\tau) \rangle d\tau - \\ &- \langle \int_0^t e(t, \tau) n(y, \dot{y}, \dots) d\tau \rangle. \end{aligned}$$

Let's examine the complicating last term for an example, say $n = \beta y^2$. The term becomes

$$- \langle \int_0^t e(t, \tau) \beta y^2(\tau) d\tau \rangle = - \langle L^{-1} \beta y^2 \rangle$$

$L^{-1} \beta y^2$ adds to the previous linear solution the terms

$$L^{-1} \beta (y_0^2 + y_1^2 + y_2^2 + \dots - 2y_0 y_1 + 2y_0 y_2 - 2y_1 y_2 + 2y_1 y_3 - \dots).$$

Thus the nonlinear term adds to $\langle y \rangle$ a series

$$\begin{aligned} &- \langle \int_0^t e(t, \tau) \beta y_0^2(\tau) d\tau \rangle + \dots = \\ &= - \int_0^t e(t, \tau) \langle \beta(\tau) \rangle \int_0^{\tau} \int_0^{\gamma} e(\tau, \gamma) e(\gamma, \sigma) \langle x(\gamma) x(\sigma) \rangle d\sigma d\gamma d\tau + \dots \end{aligned}$$

where the separation of averages takes place as before.

Now consider $R_y(t_1, t_2) = \langle y(t_1) y(t_2) \rangle$

$$\begin{aligned} R_y(t_1, t_2) &= \langle [F(t_1, \omega) - \int_0^{t_1} \Gamma(t_1, \tau_1, \omega) F(\tau_1, \omega) d\tau_1 - \\ &- \int_0^{t_1} e(t_1, \tau_1) n(y(\tau_1), \dot{y}(\tau_1), \dots) d\tau_1] \cdot [F(t_2, \omega) - \\ &- \int_0^{t_2} \Gamma(t_2, \tau_2, \omega) F(\tau_2, \omega) d\tau_2 - \end{aligned}$$

$$- \int_0^{t_2} e(t_2, \tau_2) n(y(\tau_2), y(\tau_2), \dots) d\tau_2]^* \rangle.$$

For the linear case this gave

$$\begin{aligned} R_y(t_1, t_2) = R_{FF}(t_1, t_2) - \int_0^{t_1} \langle \Gamma(t_1, \tau_1, \omega) R_{FF}(\tau_1, t_2) d\tau_1 - \\ - \int_0^{t_1} \langle \Gamma(t_2, \tau_2, \omega) \rangle R_{FF}(t_1, \tau_2) d\tau_2 + \\ + \int_0^{t_1} \int_0^{t_2} \langle \Gamma(t_1, \tau_1, \omega) \Gamma(t_2, \tau_2, \omega) \rangle R_{FF}(\tau_1, \tau_2) d\tau_1 d\tau_2 \end{aligned}$$

which, as earlier work has shown [4] reduces to the result for perturbation theory *where perturbation theory is applicable*. In the nonlinear case, we obtain nine terms, five of which involve n .

$$\begin{aligned} R_y(t_1, t_2) = \langle F(t_1) \dot{F}(t_2) \rangle - \int_0^{t_1} \langle \Gamma(t_2, \tau_2) \rangle \langle F(t_1) \dot{F}(\tau_2) \rangle d\tau_2 - \\ - \int_0^{t_1} \langle \Gamma(t_1, \tau_1) \rangle \langle F(\tau_1) \dot{F}(t_2) \rangle d\tau_1 - \\ - \int_0^{t_1} \int_0^{t_2} \langle \Gamma(t_1, \tau_1) \Gamma(t_2, \tau_2) \rangle \langle F(\tau_1) \dot{F}(\tau_2) \rangle d\tau_1 d\tau_2 - \\ - \langle \int_0^{t_2} F(t_1) \dot{e}(t_2, \tau_2) n(y(\tau_2), \dots) d\tau_2 \rangle - \\ - \langle \int_0^{t_1} \dot{E}(t_2) e(t_1, \tau_1) n(y(\tau_1), \dots) d\tau_1 \rangle + \\ + \langle \int_0^{t_1} \Gamma(t_1, \tau_1) F(\tau_1) d\tau_1 \int_0^{t_2} \dot{e}(t_2, \tau_2) n(y(\tau_2), \dots) d\tau_2 \rangle + \\ + \langle \int_0^{t_1} e(t_1, \tau_1) n(y(\tau_1), \dots) d\tau_1 \int_0^{t_2} \Gamma(t_2, \tau_2) \dot{F}(\tau_2) d\tau_2 \rangle + \\ + \int_0^{t_1} e(t_1, \tau_1) n(y(\tau_1), \dots) d\tau_1 \int_0^{t_2} \dot{e}(t_2, \tau_2) n(y(\tau_2), \dots) d\tau_2 \rangle. \end{aligned}$$

Thus the iterative process will, in the nonlinear case, add the terms (before averaging)

$$\begin{aligned}
 & - \int_0^{t_2} F(t_1) e^*(t_2, \tau_2) \dot{n}(y_0(\tau_2), \dots) d\tau_2 - \\
 & - \int_0^{t_1} \dot{F}(t_2) e(t_1, \tau_1) n(y_0(\tau_1), \dots) d\tau_1 + \\
 & + \int_0^{t_1} \Gamma(t_1, \tau_1) F(\tau_1) d\tau_1 \int_0^{t_2} e^*(t_2, \tau_2) \dot{n}(y_0(\tau_2), \dots) d\tau_2 + \\
 & + \int_0^{t_2} \Gamma(t_2, \tau_2) \dot{F}(\tau_2) d\tau_2 \int_0^{t_1} e(t_1, \tau_1) n(y_0(\tau_1), \dots) d\tau_1 + \\
 & + \int_0^{t_1} e(t_1, \tau_1) n(y_0(\tau_1), \dots) d\tau_1 \int_0^{t_2} e^*(t_2, \tau_2) \dot{n}(y_0(\tau_2), \dots) d\tau_2
 \end{aligned}$$

+ higher terms involving $y_1, y_2, \dots$.

Suppose as an example again $n(y(\tau)) = \beta y^2(\tau)$ we will obtain (after averaging them) the additional terms to $R_y(t_1, t_2)$ of

$$\begin{aligned}
 & - \int_0^{t_2} F(t_1) e^*(t_2, \tau_2) \beta y_0^2(\tau_2) d\tau_2 - \\
 & - \int_0^{t_1} \dot{F}(t_2) e(t_1, \tau_1) \beta y_0^2(\tau_1) d\tau_1 + \\
 & + \int_0^{t_1} \Gamma(t_1, \tau_1) F(\tau_1) d\tau_1 \int_0^{t_2} e^*(t_2, \tau_2) \beta y_0^2(\tau_2) d\tau_2 + \\
 & + \int_0^{t_2} \Gamma(t_2, \tau_2) \dot{F}(\tau_2) d\tau_2 \int_0^{t_1} e(t_1, \tau_1) \beta y_0^2(\tau_1) d\tau_1 + \\
 & + \int_0^{t_1} e(t_1, \tau_1) \beta y_0^2(\tau_1) d\tau_1 \int_0^{t_2} e^*(t_2, \tau_2) \beta y_0^2(\tau_2) d\tau_2
 \end{aligned}$$

and the higher terms of the iteration [1]. We see averages will all separate.

In the linear case, $\langle y \rangle$ was a two-fold integral of $y(t, \omega)$ over the measure spaces of the x process and the coefficients α_i . In the nonlinear case, we have an additional integration because of the $b_\mu(t, \omega)$.

However, because of the statistical independence of coefficients and input and the iteration back to $F(t)$, the separations will occur as before the n term gives us an integral over the measure space of the b_n as discussed earlier [1]. A number of special cases have been considered and will appear elsewhere; the complete theory with applications will appear in a book edited by Richard Bellman.

University of Georgia
Athens, Georgia, U. S. A.

Received 8.X.1975

Գ. Ա. ԱԴՈՄՅԱՆ. Ոչ գծային ստոխաստիկ դիֆերենցիալ հավասարումների առաջին և երկրորդ կարգի $\langle y \rangle$ և $R_y(t_1, t_2)$ ստատիստիկաների ստացումը (ամփոփում)

Հոգվածում ինտերացիայի մեթոդով ստացված են երկու կորելացիայի ֆունկցիաներ $R_y(t_1, t_2)$ և $\langle y \rangle$ միջին լուծումը ընդհանուր ոչ գծային ստոխաստիկ դիֆերենցիալ հավասարման համար:

Այդ մեթոդը հանդիսանում է Ադոմյանի (գծային հավասարման համար) մեթոդի ընդհանուրացումը:

Г. А. АДОМЯН. Получение статистик первого и второго порядка, $\langle y \rangle$ и $R_y(t_1, t_2)$ для нелинейных стохастических дифференциальных уравнений (резюме)

Для общего нелинейного стохастического дифференциального уравнения среднее решение $\langle y \rangle$ и две точечные функции корреляции $R_y(t_1, t_2)$ получены методом итерации, обобщающего метод Адомяна для линейного стохастического уравнения.

REFERENCES

1. G. Adomian. Nonlinear Stochastic Differential Equations, Jour. of Mathematical Analysis and Applications, in publication.
2. G. Adomian. Stochastic Systems Modelling, Third International Congress of Cybernetics and Systems, Bucharest, Romania, 25—29 August, 1975 (Proceedings to be published in book form).
3. G. Adomian. Random Operator Equations in Mathematical Physics I, Journal of Mathematical Physics, vol. 11, no. 3, pp. 1069—1084, March, 1970.
4. G. Adomian. Random Operator Equations in Mathematical Physics II, Journal of Mathematical Physics, vol. 12, no. 9, pp. 1944—48, September, 1971.
5. G. Adomian. The Closure Approximation in the Hierarchy Equations, Journal of Statistical Physics, vol. 3, no. 2, Summer, 1971.
6. L. H. Sibul. Application of Linear Stochastic Operator Theory, Ph. D. Dissertation, Pennsylvania State University, 1968.